

# Computing charm and bottom quark masses in lattice QCD



SFB/TR 9

Rainer Sommer

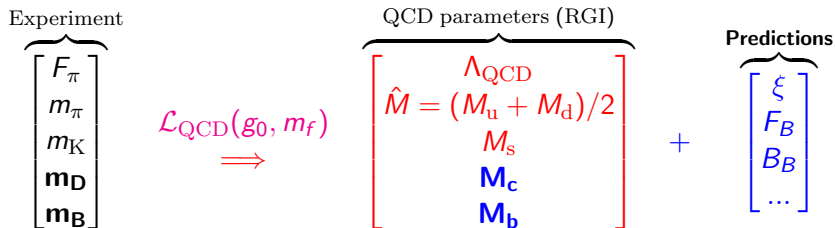
DESY, Zeuthen



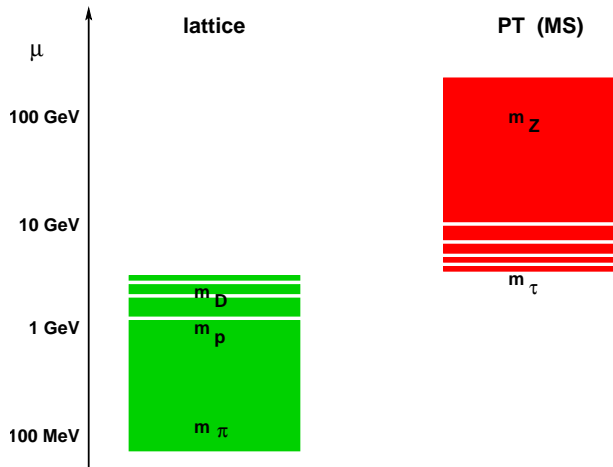
SFB workshop, Aachen February 2007

Introduction  
Charm  
Bottom  
Challenges

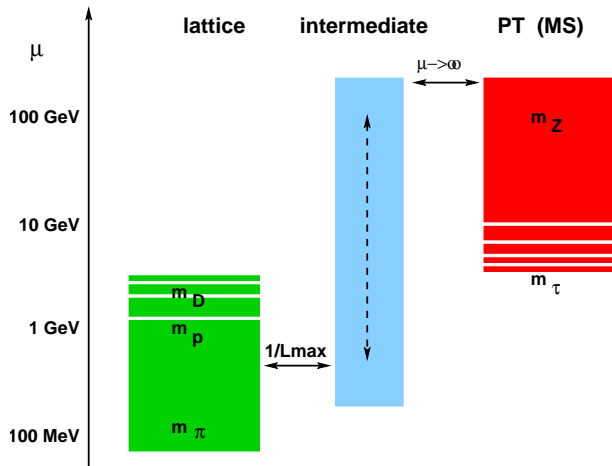
# What do we want



# Scale problem and strategy



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- ▶ intermediate: Schrödinger functional scheme

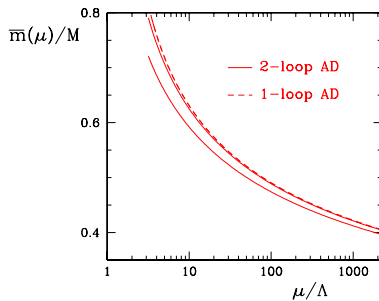
## RGI quark masses

$$M = \lim_{\mu \rightarrow \infty} \bar{m}(\mu) [2b_0 \bar{g}(\mu)^2]^{-d_0/2b_0}$$

$$\mu \frac{\partial \bar{m}}{\partial \mu} = \tau(\bar{g})$$

$$\tau(\bar{g}) \stackrel{\bar{g} \rightarrow 0}{\sim} -\bar{g}^2 \{d_0 + \bar{g}^2 d_1 + \dots\},$$

$$d_0 = \frac{8}{(4\pi)^2}$$



- $M$  scheme & scale independent:

$$\frac{\bar{m}_1(\mu)}{\bar{m}_2(\mu)} = 1 + O(\alpha(\mu))$$

$$\xrightarrow{\mu \rightarrow \infty} 1 \quad \rightarrow \quad M_1 = M_2$$

# Running mass, definition

- ▶ from the PCAC relation:

$$\begin{aligned} A_{\mu}^{SC} &= \bar{s} \gamma_{\mu} \gamma_5 c, \quad P^{SC} = \bar{s} \gamma_5 c \\ \partial_{\mu} A_{\mu}^{SC} &= (m_c + m_s) P^{SC} \end{aligned}$$

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- ▶ precisely

$$g_0 \leftrightarrow a$$

$$Z_A(g_0) \langle D_s(p=0) | A_\mu^{SC} | 0 \rangle = (\bar{m}_c(\mu) + \bar{m}_s(\mu)) Z_P(\mu, g_0) \langle D_s(p=0) | P^{SC} | 0 \rangle$$

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- ▶ or

$$\underbrace{\bar{m}_c(\mu)}_{\text{renormalized, running}} = \frac{Z_A(g_0)}{Z_P(\mu, g_0)} \underbrace{m_c(g_0)}_{\text{bare, PCAC}}$$

scheme dependence from  $Z_P$



# The basic equation

$$\bar{m}_c(\mu) = \frac{Z_A(g_0)}{Z_P(\mu, g_0)} \underbrace{m_c(g_0)}_{\text{bare, PCAC}}$$

$$M_c = \frac{M}{\underbrace{\bar{m}(\mu)}} \bar{m}_c(\mu)$$

talk by FK

$$M_c = Z_M(g_0) \underbrace{m_c(g_0)}_{\text{bare, PCAC}}, \quad Z_M(g_0) = \frac{M}{\bar{m}(\mu)} \frac{Z_A(g_0)}{Z_P(\mu, g_0)}$$

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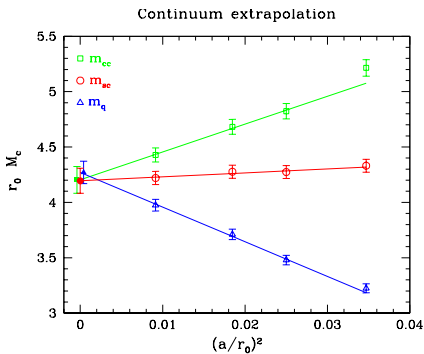
- ▶ physics input: bare charm mass in the Lagrangian s.t.  $m_D/F_\pi = \text{experiment}$
- ▶  $Z_M(g_0)$  is known **non-perturbatively**  
for Wilson-type LQCD,  $N_f = 0, 2$

# Charm quark mass

- remove  $O(am_c)$  effects non-perturbatively [ALPHA Collaboration 2000]

quenched [Sint & Rolf]

charm **just** doable

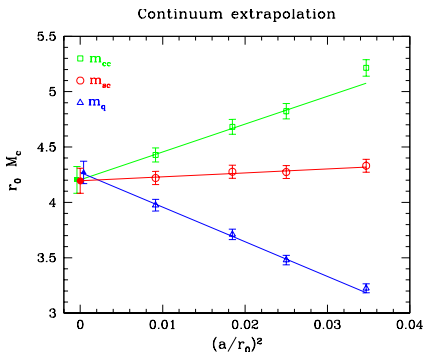


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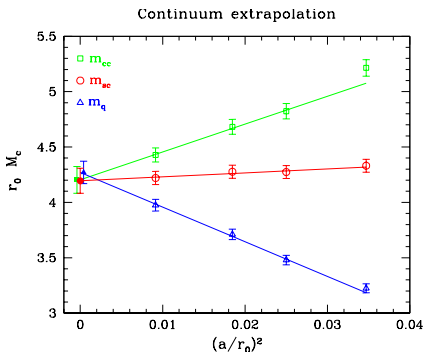
- $N_f = 0$ ,  $r_0 = 0.5$  fm:  $M_c = 1654(45)\text{MeV} \rightarrow \bar{m}_c^{\overline{\text{MS}}}(\bar{m}_c) = 1301(34)\text{MeV}$

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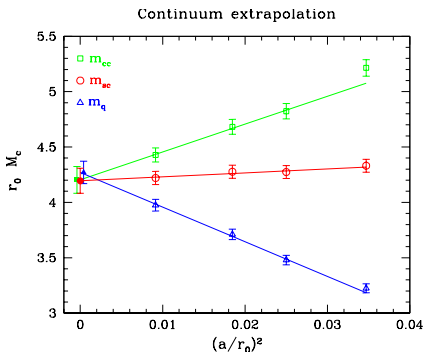
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- $aM_b \approx 4aM_c$  !  
for b-quarks  $a \rightarrow 0$  can't be controlled in this way  $\rightarrow$  eff. theory

# Effective theory for b-quark in heavy-light systems: HQET

[Caswell & Lepage ... Eichten & Hill ... 

$$\bar{b}(x)[\gamma_\mu D_\mu + m_b]b(x) \rightarrow$$

$$\mathcal{L}_{\text{HQET}} = a^4 \sum_x \{ \bar{\psi}_h(x)[D_0 + \delta m]\psi_h(x)$$

$$+ \underbrace{\omega_{\text{spin}}}_{\sim 1/2m_b} \bar{\psi}_h(-\boldsymbol{\sigma} \cdot \mathbf{B})\psi_h + \underbrace{\omega_{\text{kin}}}_{\sim 1/2m_b} \bar{\psi}_h(-\frac{1}{2}\mathbf{D}^2)\psi_h + \dots \}$$

$$P_+ \psi_h = \psi_h, \quad \bar{\psi}_h P_+ = \bar{\psi}_h, \quad P_+ = \frac{1}{2}(1 + \gamma_0)$$



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non-trivial matching problem

$$\omega_{\text{kin}}, \omega_{\text{spin}}, \delta m \Leftrightarrow M_b$$

# Mass renormalization in HQET

► static:  $\bar{\psi}_h(x)[D_0 + \delta m]\psi_h(x)$

$$m_b^{\overline{\text{MS}}} = Z^{\overline{\text{MS}}} [m_{\text{bare}} + \delta m]$$

$$m_{\text{bare}} = \underbrace{m_B}_{\text{exp.}} - \underbrace{E_{\text{stat}}}_{\text{“binding energy”}}$$

↑

$$= \frac{e(g_0)}{a} \sim \exp(1/(2b_0g_0^2)) [1 + e_1g_0^2 + \dots]$$

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no continuum limit
- ▶ need non-perturbative  $e(g_0)$   
more generally: non-perturbative matching (renormalisation) of HQET

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more generally: non-perturbative matching (renormalisation) of HQET
- ▶  $1/m$  corrections make it worse:  $a^{-1} \rightarrow a^{-2}$

# Non-perturbative matching

►  $\omega_{\text{kin}}, \omega_{\text{spin}}, \delta m + m_{\text{bare}}$  from QCD

[Heitger & S.]

$$\delta m + m_{\text{bare}} \leftrightarrow M_b$$

# Non-perturbative matching

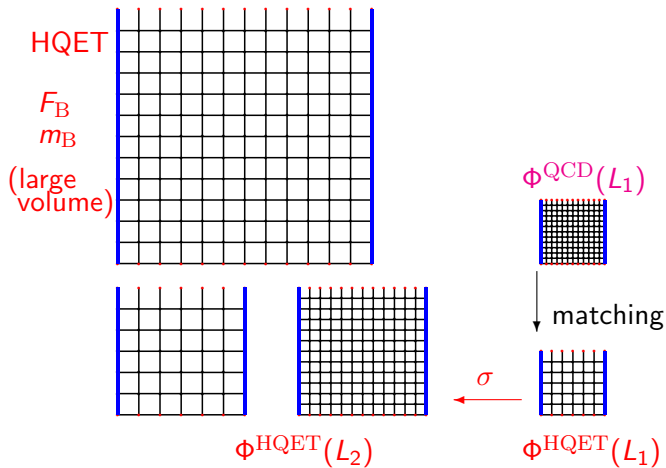
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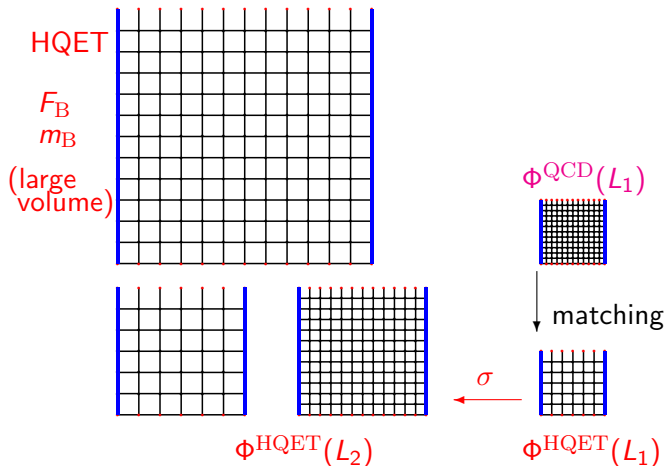
$$\delta m + m_{\text{bare}} \leftrightarrow M_{\text{b}}$$

- ▶  $\Phi_i^{\text{QCD}}(L) = \Phi_i^{\text{HQET}}(L), \quad i = 1, 2, 3$

# The full strategy



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- continuum limit can be taken in all steps



# Equation for $M_b$ (static)

► pert:  $m_B = [Z^{\overline{\text{MS}}}]^{-1} m_b^{\overline{\text{MS}}} + \underbrace{E_{\text{stat}}}_{\text{"binding energy"}} + a^{-1}[1 + e_1 g_0^2 + \dots]$

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▶ NP:

$$m_B = \underbrace{E_{\text{stat}} - E_{\text{stat}}^{\text{sub}}}_{E_{\text{stat}} - E_{\text{stat}}(L_2) + E_{\text{stat}}(L_2) - E_{\text{stat}}(L_1)} + \underbrace{E_{\text{stat}}^{\text{sub}}}_{E(L_1, M_b)} + \underbrace{E_{\text{stat}}^{\text{sub}}}_{\text{QCD}}$$

HQET

$$L_1 \approx 0.4 \text{ fm}, \quad L_2 = 2L_1$$

## Example for continuum extrapolation (quenched)

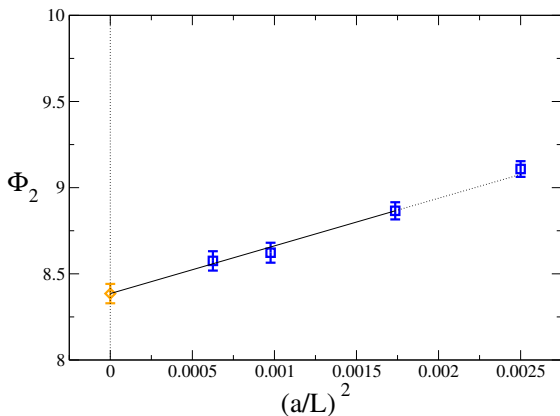
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3 values of  $M_b$   
example:

( $M_b$  from bare  
mass as before:  
 $Z_M$ )

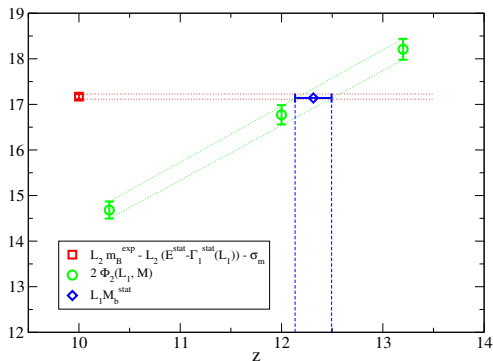


# Result

in static approximation

express everything in units of  $r_0 \approx 0.5$  fm and solve:

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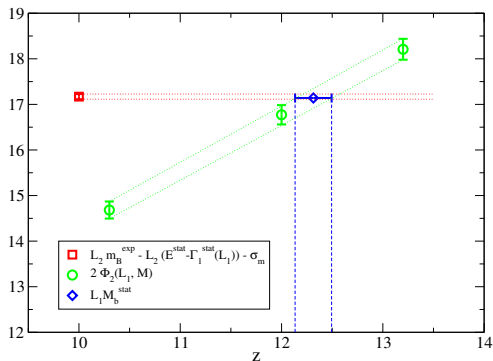


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with  $r_0 = 0.5$  fm

$$M_b^{\text{stat}} = \begin{cases} 6771 \pm 99 \text{ MeV (HYP2)} \\ 6757 \pm 99 \text{ MeV (HYP1)} \end{cases}$$

and obtain the slope

$$S = \frac{1}{L_1} \frac{\partial \Phi^{\text{QCD}}(L_1, M)}{\partial M} = 0.61(5)$$

error is dominated by the one on  $Z_M$ .

# Results for different matching observables

- ▶ included  $1/m_b$  corrections

[Della Morte, Garron, Papinutto, S.]

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		$\theta_1 = 0$ $\theta_2 = 1/2$	$\theta_1 = 1/2$ $\theta_2 = 1$	$\theta_1 = 1$ $\theta_2 = 0$
		Main strategy		
0	17.25(20)	17.12(22)	17.12(22)	17.12(22)
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0	17.05(25)	17.25(28)	17.23(27)	17.24(27)
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- ▶ related strategy, HQET only at lowest order,  $1/m_b$  corrections from QCD interpolations  $\bar{m}_b^{\overline{\text{MS}}}(\bar{m}_b) = 4.421(67)\text{GeV}$

[Guazzini, S., Tantalò]

talk by Damiano Guazzini

# Challenges, todos

- ▶ **error** for  $M_b$  dominated by  $Z_M(g_0)$  in relativistic QCD  
→ **reduce by a factor 2 ?!**  
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- ▶  $N_f = 3 \dots$

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- ▶ no assumptions about (local) duality, “convergence” of the (divergent) pert. expansion, renormalon subtraction, ...