

Status of the Strange Quark Mass Determination

Alexander Khodjamirian



Mini-Workshop on α_s and quark masses, Aachen, 21.02.07

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 $\sim 25\%$ accuracy ,
compared with $\sim 10\%$ for $\overline{m}_c(\overline{m}_c)$ and $\sim 2\%$ for $\overline{m}_b(\overline{m}_b)$

- ChPT yields the relations:

$$R = \frac{m_s}{\hat{m}} = 24.4 \pm 1.5, \quad Q^2 = \frac{m_s^2 - \hat{m}^2}{m_d^2 - m_u^2} = (22.7 \pm 0.8)^2$$

[Leutwyler, 1996]

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$$\hat{m} = \frac{1}{2}(m_u + m_d)$$

determine m_s , obtain $m_{u,d}$ “for free “:

$$m_d = \frac{m_s}{R} \left(1 + \frac{R-1}{4Q^2} \right)$$
$$m_u = \frac{m_s}{R} \left(1 - \frac{R-1}{4Q^2} \right)$$

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 - * inclusive $\tau \rightarrow s\bar{u}\nu_\tau$ decays
 - * Finite-energy sum rules (FESR)

Lattice QCD results (≥ 2005)

n_f	$\overline{m}_s(2\text{GeV})[\text{MeV}]$	Ref.
2	97 ± 22	[ALPHA Collab.] Della Morte et al Nucl.Phys. B (2005)
2	$101 \pm 8_{-0}^{+25}$	[SPQ _{CD} R Collab.] Becirevic et al. Proc.Lat2005
2	$115 \pm 2 \pm 3 \pm 6$	[QCDSF-UKQCD Collab.] M. Göckeler et al. Proc.Lat2006
3	$87 \pm 4 \pm 4$	[HPQCD Collab.] Mason et al. PRD (2006)
3	$91.8 \pm 3.9_{-4.1}^{+6.8}$	[CP-PACS & JLQCD Collab.], Ishikawa et al., Proc.Lat2006
3	$90 \pm 5 \pm 4$	[MILC Collab.] Bernard et al.Proc.Lat2006
2,3	95 ± 20	<i>lattice average</i> , C. T. Sachranda, A. V. Manohar PDG review on m_q (2006)

Employing quark-current correlators

$$\Pi(q^2) = i \int d^4x e^{iq \cdot x} \langle 0 | T \{ j^{(s)}(x) j^{(s)\dagger}(0) \} | 0 \rangle$$

$$j^{(s)} = (\bar{s}q)_{V,A,S,P,\dots}$$

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* OPE at $Q^2 = -q^2 \gg \Lambda_{QCD}^2$
(PQCD \oplus condensates)

$$\Pi(Q^2) = \sum_{d=1,2,3,\dots} \frac{C_d(\alpha_s, m_s, \dots)}{(Q^2)^d}$$

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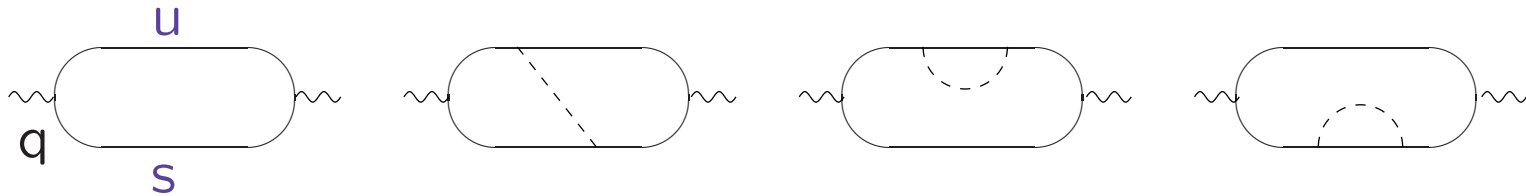
$$\Pi(Q^2) = \sum_{d=1,2,3,\dots} \frac{C_d(\alpha_s, m_s, \dots)}{(Q^2)^d} = \frac{1}{\pi} \int_0^{\infty} ds \frac{\text{Im}\Pi(s)}{(s + Q^2)} \quad (+subtr.)$$

* unitarity:

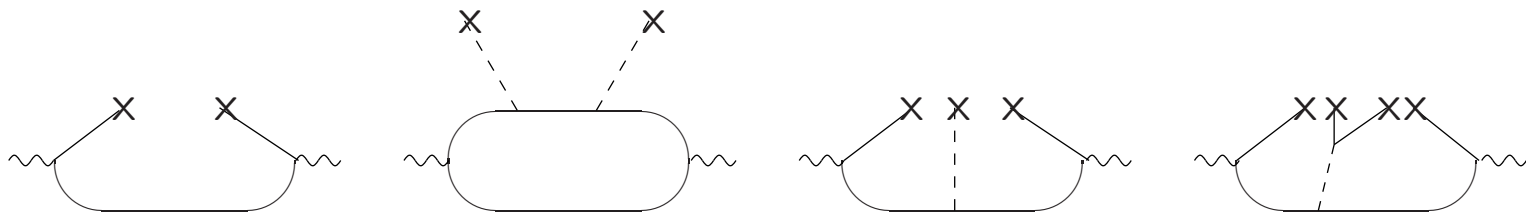
$$\frac{1}{\pi} \text{Im}\Pi(s) = \sum_{h_s} \langle 0 | j^{(s)} | h_s \rangle \langle h_s | j^{(s)} | 0 \rangle \leftarrow \text{hadronic observables}$$

e.g., $j_P^{(s)} = \bar{s}\gamma_5 q$, $h_s = K, K2\pi, K^*\pi, \dots$

OPE: perturbative and condensate diagrams



$$\oplus O(\alpha_s^2) \oplus O(\alpha_s^3) \oplus O(\alpha_s^4)$$



$\langle \bar{q}q \rangle$

$\langle GG \rangle$

$\langle \bar{q}Gq \rangle$

$\langle \bar{q}q \rangle^2$

$$\oplus O(\alpha_s)$$

Current status of OPE

- Correlators with vector and axial currents (employed in τ -decays):
calculated to order α_s^3 [Baikov, Chetyrkin, Kühn, PRL(2005)]
- Correlators with scalar (pseudoscalar) currents:

$$j_{S(P)}^{(s)} = \partial^\mu \bar{s} \gamma_\mu (\gamma_5) q = (m_s \mp m_q) \bar{s} (\gamma_5) q, \quad (q = u, d)$$

$$\Pi(Q^2) = (m_s \mp m_q)^2 \tilde{\Pi}(Q^2)$$

calculated to order α_s^4 [Baikov, Chetyrkin, Kühn, PRL(2006)]

OPE for the pseudoscalar correlator

expansion in $1/(Q)^{d+2}$, $d = 0, 2, 4, 6$

$$\begin{aligned} [\Pi^{(P)''}(Q^2)]_{OPE} &= \frac{3(m_s + m_u)^2}{8\pi^2 Q^2} \left\{ 1 + \sum_{i=1}^4 C_{0,i} \left(\frac{\alpha_s}{\pi}\right)^i \right. \\ &\quad \left. - 2\frac{m_s^2}{Q^2} \left(1 + \sum_{i=1,2} C_{2,i} \left(\frac{\alpha_s}{\pi}\right)^i \right) + \frac{\{d=4\}}{Q^4} + \frac{\{d=6\}}{Q^6} \right\} \end{aligned}$$

$$\{d=4\} \sim \{m_s \langle \bar{q}q \rangle, \langle G^2 \rangle, O(m_s^4)\} (1 \oplus O(\alpha_s))$$

$$\{d=6\} \sim m_s \langle \bar{q}Gq \rangle, \langle \bar{q}q \rangle^2$$

Coefficients multiplying $(\alpha_s/\pi)^n$ in $d = 0$ part: ($l_Q = \log Q^2/\mu^2$)

$$C_{0,1} = \frac{11}{3} - 2l_Q, \quad C_{0,2} = \frac{5071}{144} - \frac{35}{2} \zeta_3 - \frac{139}{6} l_Q + \frac{17}{4} l_Q^2,$$

$$C_{0,3} = \frac{1995097}{5184} - \frac{\pi^4}{36} - \frac{65869}{216} \zeta_3 + \frac{715}{12} \zeta_5 - \frac{2720}{9} l_Q + \frac{475}{4} \zeta_3 l_Q + \frac{695}{8} l_Q^2 - \frac{221}{24} l_Q^3,$$

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new:

$$\begin{aligned} C_{0,4} = & \frac{2361295759}{497664} - \frac{2915}{10368} \pi^4 - \frac{25214831}{5184} \zeta_3 + \frac{192155}{216} \zeta_3^2 + \frac{59875}{108} \zeta_5 - \frac{625}{48} \zeta_6 \\ & - \frac{52255}{256} \zeta_7 + l_Q \left[-\frac{43647875}{10368} + \frac{1}{18} \pi^4 + \frac{864685}{288} \zeta_3 - \frac{24025}{48} \zeta_5 \right] \\ & + l_Q^2 \left[\frac{1778273}{1152} - \frac{16785}{32} \zeta_3 \right] + l_Q^3 \left[-\frac{79333}{288} \right] + l_Q^4 \left[\frac{7735}{384} \right], \end{aligned}$$

Hierarchy in α_s and d

Relative contributions to $[\Pi^{(P)''}(M^2)]_{OPE}$ (after Borel transf. $Q^2 \rightarrow M^2$)

$$r_n^{(d)}(M^2) = \frac{\{\Pi^{(P)''}(M^2)\}_{O(\alpha_s^n)}^{(d)}}{\Pi^{(P)''}(M^2)}$$

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$$r_n^{(d)}(M^2) = \frac{\{\Pi^{(P)''}(M^2)\}_{O(\alpha_s^n)}^{(d)}}{\Pi^{(P)''}(M^2)}$$

$$r_n^{(d=0,2)}(2.5 \text{ GeV}^2) = 52.4\%, 28.3\%, 14.4\%, 4.0\%, -0.3\%, n = 0, 1, 2, 3, 4,$$

$$r^{(d=4,6)}(2.5 \text{ GeV}^2) = 1.2\%.$$

Bounds for m_s

- positivity of the spectral function e.g., in the pseudoscalar channel:

$$\frac{1}{\pi} \text{Im}\Pi^{(P)}(s) = f_K^2 m_K^4 \delta(m_K^2 - s) + \{\geq 0\}$$

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- double-subtracted dispersion relation $\Pi^{(P)''}(Q^2)$ yields the linear bound:

$$[m_s(Q^2) + m_q(Q^2)]^2 \geq \frac{32\pi^2}{3Q^4} \frac{f_K^2 m_K^4}{(1 + m_K^2/Q^2)^3} \left(\frac{1}{1 + O(\alpha_s) + \dots} \right)$$

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- power moments: $\mathcal{F}_n(Q^2) = c_n(Q^2)^{n+1} \left(\frac{d}{dQ^2} \right)^{n+2} \Pi''(Q^2)$

\Rightarrow the quadratic bound, combining $\mathcal{F}_{0,1,2}$,

[Lellouch, de Rafael, Taron (1997)]

● recent $O(\alpha_s^4)$ improvement [Baikov, Chetyrkin, Kühn, PRL(2006)]

* consistent α_s -expansion of the ratio of moments

* good convergence of perturb. series

Bound	$[\overline{m}_s + \overline{m}_u](2\text{GeV})$	precision
linear	$> 72 \text{ MeV}$	$O(\alpha_s^4)$
quadratic	$> 77 \text{ MeV}$	$O(\alpha_s^4)$

QCD sum rule (pseudoscalar, Borel)

Dispersion relation for $\Pi^{(P)}(q^2)$ at large Q^2 (doubly differentiated):

$$[\Pi^{(P)''}(q^2)]_{OPE} = 2 \int_0^{\infty} ds \frac{\rho^{(P)}(s)}{(s - q^2)^3} .$$

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$$\rho^{(P)}(s) = \sum_K \langle 0 | j_P^{(s)} | K(q) \rangle \langle K(q) | j_P^{(s)} | 0 \rangle$$

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Borel transformation:

$$M^4 [\Pi^{(P)''}(M^2)]_{OPE} = \int_0^{s_0} ds e^{-s/M^2} \rho_{hadr}^{(P)}(s) + \int_{s_0}^{\infty} ds e^{-s/M^2} \rho_{OPE}^{(P)}(s).$$

The hadronic part

- $K, K2\pi, K^*\pi, \rho K, \dots \rightarrow$ three-resonance ansatz $\{K, K_1(1460), K_2(1830)\}$
 $m_{K_1}=1460$ MeV, $\Gamma_{K_1} = 260$ MeV; $m_{K_2}=1830$ MeV, $\Gamma_{K_2}=250$ MeV [PDG]

$$\rho_{had}^{(P)}(s) = f_K^2 m_K^4 \delta(m_K^2 - s) + \sum_{i=1,2} f_{K_i}^2 m_{K_i}^4 \frac{1}{\pi} \left(\frac{\Gamma_{K_i} m_{K_i}}{(s - m_{K_i}^2)^2 + (\Gamma_{K_i} m_{K_i})^2} \right)$$

- decay constants: $\langle 0 | j_5^{(s)} | K(q) \rangle = f_K m_K^2$, $f_K = 159.8$ MeV ,
 $f_{K_1} = \sqrt{2}(22.9 \pm 2.4)$ MeV, $f_{K_2} = \sqrt{2}(14.5 \pm 1.5)$ MeV, $s_0 = 4.5 \pm 0.5$ GeV²

fitted from pseudoscalar sum rules and FESR with $O(\alpha_s^3)$ accuracy

[Kambor, Maltman , PRD 2002] (small f_{K_i} consistent with PCAC/ChPT)

$O(\alpha_s^4)$ improvement of the pseudoscalar sum rule

[Chetyrkin, A.K. , EJP (2006)]

- inputs: $\alpha_s(m_Z)$, $(\alpha_s(m_\tau))$, condensates

$$\overline{m}_u(2 \text{ GeV}) = (1.5 - 5.0) \text{ MeV},$$

$$\overline{m}_d(2 \text{ GeV}) = (5.0 - 9.0) \text{ MeV (non-lattice) intervals from [PDG]}$$

- Borel mass interval: $2 < M^2 < 3 \text{ GeV}^2$,
(to avoid instantons and to minimize the contribution of excited states)

the scale $\mu = M$, varying $0.5 < \mu^2/M^2 < 2.0$

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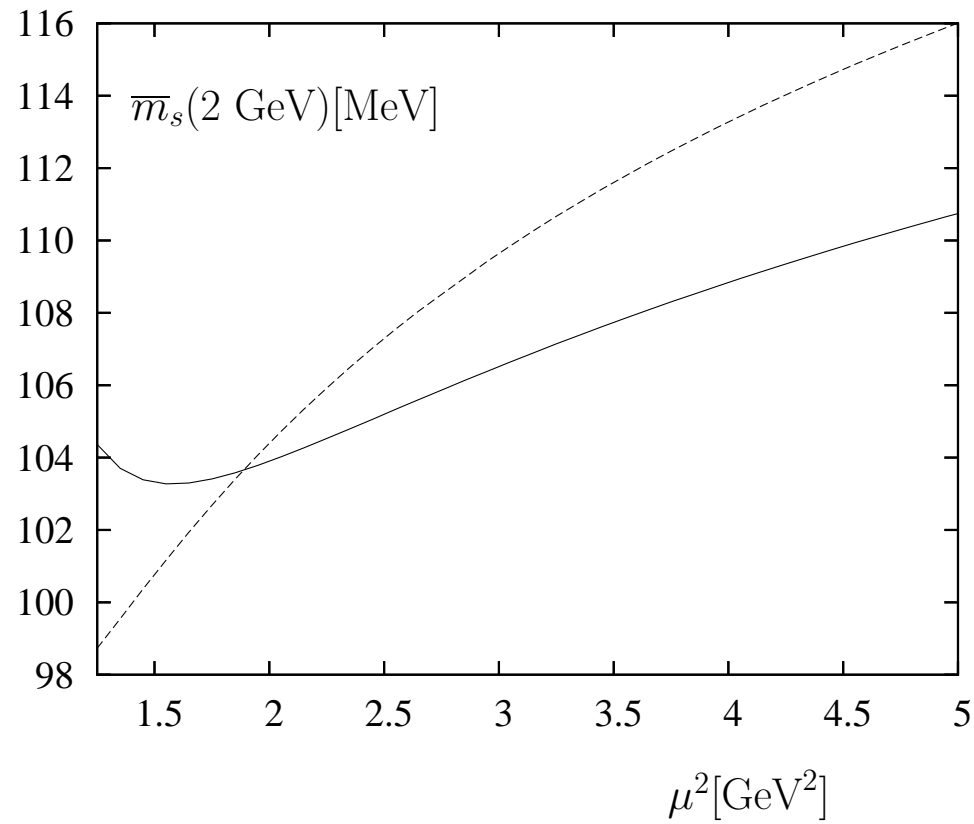
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- the result:
$$\overline{m}_s(2 \text{ GeV}) = \left(105 \pm 6 \Big|_{OPE} \pm 7 \Big|_{hadr} \right) \text{ MeV},$$

$\simeq 2 \text{ MeV}$ increase of the central value, if the $O(\alpha_s^4)$ terms are removed.

The scale dependence



solid $O(\alpha_s^4)$, dashed $O(\alpha_s^3)$

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($K_1 \rightarrow K^* \pi, K \rho, K 2\pi, \dots$),
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($K_1 \rightarrow K^* \pi, K \rho, K 2\pi, \dots$),
but: $BR \sim 10^{-5}$ -chirally and Cabibbo suppressed
- excited kaons in semileptonic charmed meson decays,
e.g., $D \rightarrow K_1(1460) l \nu_l$,
no chiral/Cabibbo suppression, but: no access to f_{K_1}
- kaon resonances in $B \rightarrow J/\psi K \pi \pi$ etc. decays, $BR \sim 10^{-4} - 10^{-3}$

other QCD sum rules

- use of the scalar correlator

- * the same OPE as in the pseudoscalar case $(m_s + m_u)^2 \rightarrow (m_s - m_u)^2$

- * a different pattern of the hadronic spectral function, $J^P = O^+$ $K\pi$ -state
a reliable knowledge of the $K \rightarrow \pi$ scalar form factor is needed

- * accuracy improved to $O(\alpha_s^4)$ [Jamin, Oller, Pich '06]

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a reliable knowledge of the $K \rightarrow \pi$ scalar form factor is needed
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- FESR (scalar, pseudoscalar) [Maltman, Kambor, '02]
vector FESR [Eidemüller, Jamin, '03]

m_s from $\tau \rightarrow s\bar{u}\nu_\tau$ decays

- * based on vector/axial-vector correlators
- * $\oint \Pi(Q^2)dQ^2$ related to $\int_0^{m_\tau^2} ds \text{Im}\Pi(s)$, where $\text{Im}\Pi(s) \sim R_\tau(s)$
- * the art of combining the moments of spectral densities [K. Maltman et al.]
- * less sensitive to m_s , serves primarily for a precise extraction of $|V_{us}|$

(see the next talk by Matthias Jamin)

Non-lattice determinations (≥ 2005)

Method	$\overline{m}_s(2 \text{ GeV})$ in MeV	Ref.
Hadronic τ decays	96_{-3-18}^{+5+16} 104 ± 28 76 ± 20 89 ± 28	Baikov, Chetyrkin, Kühn, PRL (2005) S. Narison, hep-ph/0510108 Gamiz et al, hep-ph/0612154 Maltman, Wolfe, hep-ph/0701037
Pseudoscalar Borel SR	$105 \pm 6 \pm 7$ $97.2_{-8.0}^{+11.3}$	A.K., Chetyrkin, hep-ph/0612295 Jamin, Oller, Pich hep-ph/0605095
Scalar Borel SR	$87.6_{-6.8}^{+8.8}$	Jamin, Oller, Pich hep-ph/0605095
nonlattice average	103 ± 20	C. T. Sachranda, A. V. Manohar PDG review on m_q (2006) $O(\alpha_s^4)$ not included

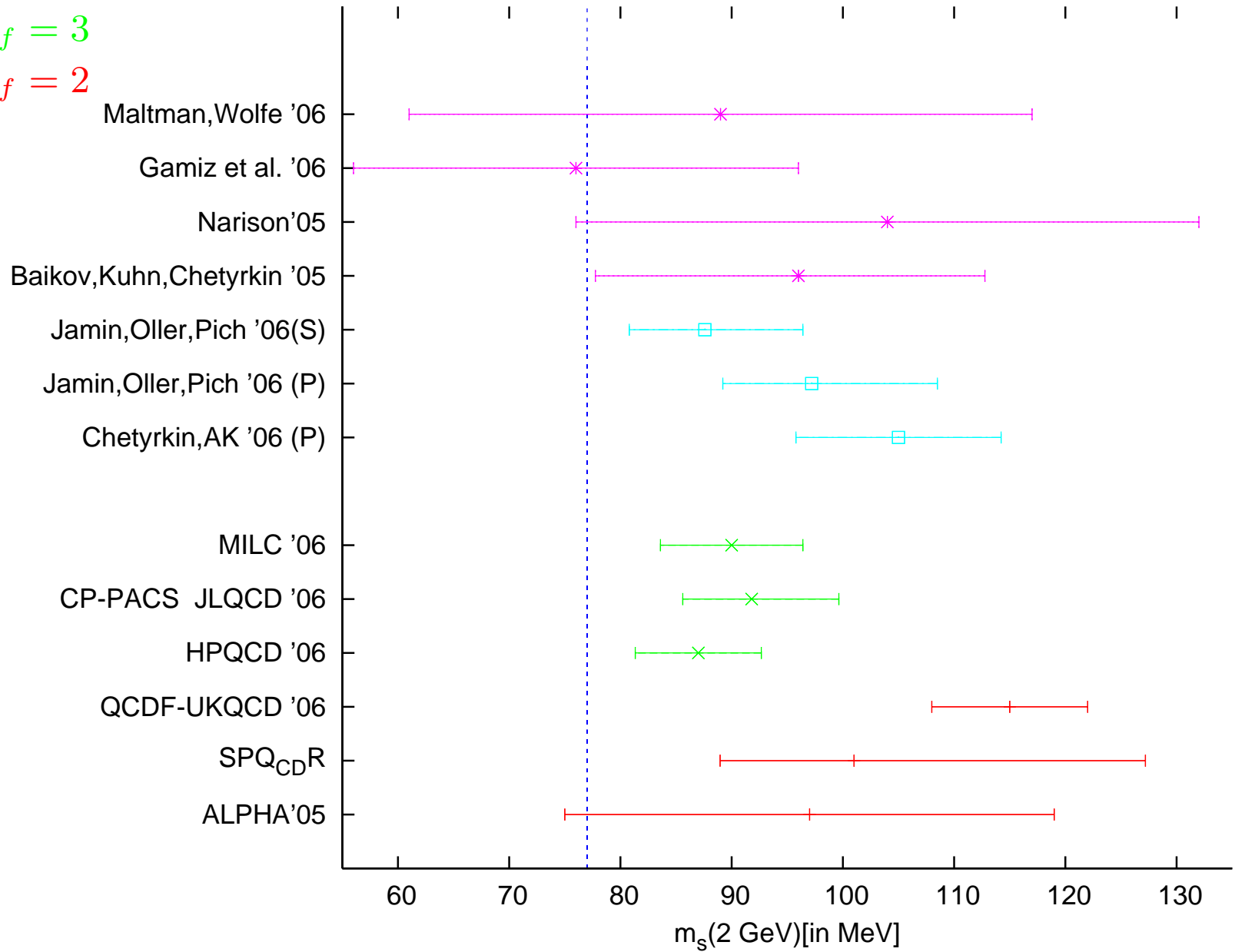
τ decays

QCD SR

LQCD $n_f = 3$

LQCD $n_f = 2$

quadr.bound, Baikov, Chetyrkin, Kühn '06

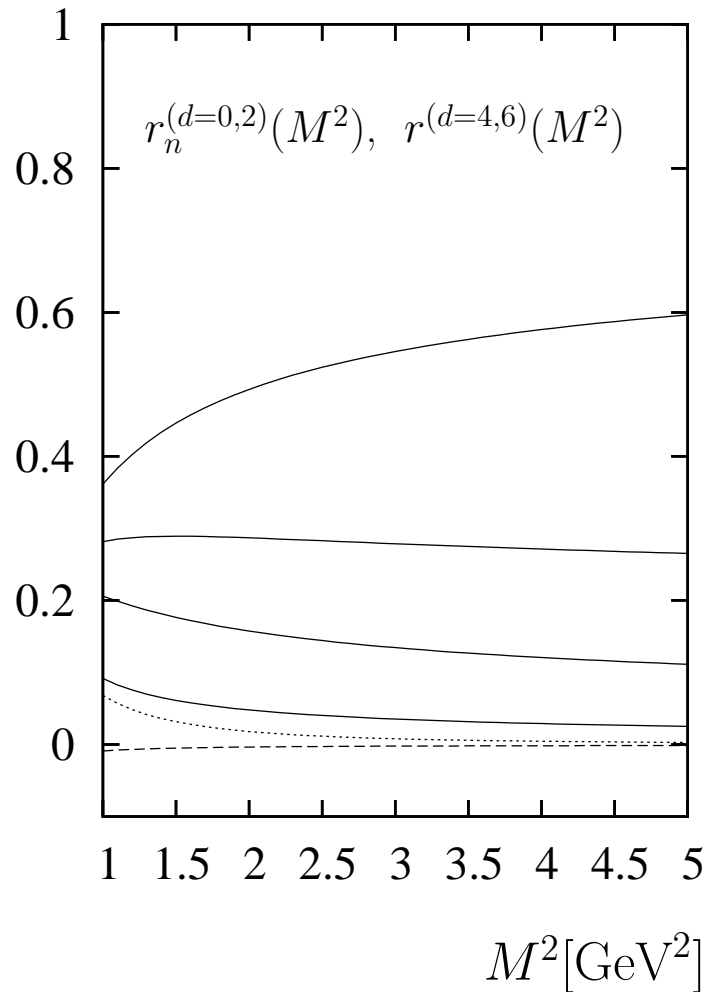


Summary

- $O(\alpha_s^4)$ term in the correlators of (pseudo)scalar $\Delta S = 1$ quark currents:
good convergence of OPE in both α_s and d
- the bound on m_s improved by including $O(\alpha_s^4)$
- QCD sum rules with $O(\alpha_s^4)$ correction are less scale-dependent than in $O(\alpha_s^3)$
- hadronic input in the sum rules can be refined with more experimental data on kaon resonances, e.g., from exclusive τ , D and B decays
- lattice vs “continuum” determinations - not only a general agreement:
retaining only $n_f = 3$ lattice and $O(\alpha_s^4)$ QCDSR,
(PDG06) interval $m_s(2\text{GeV}) = 70 - 120 \text{ MeV} \Rightarrow m_s(2\text{GeV}) = 82 - 115 \text{ MeV}$.

Backup Slides

- Relative contributions to OPE: dependence on Borel mass



solid (from up to down):
 $r_n^{(d=0,2)}$ with $n = 0, 1, 2, 3$,
dashed: $r_4^{(d=0,2)}$,
dotted: $r^{(d=4,6)}$.

Results for $\overline{m}_s(2 \text{ GeV})$, as end of 2005

Pseudoscalar Borel sum rule	$105 \pm 6 \pm 7$ 100 ± 6	This work [Maltman, Kambor '00]
Pseudoscalar FESR	100 ± 12	[Maltman, Kambor '00]
Scalar Borel sum rule	99 ± 16	[Jamin,Oller,Pich '02]
Vector FESR	139 ± 31	[Eidemuller,Jamin,Schwab '03]
Hadronic τ decays	81 ± 22	[Gamiz et al,'05]
	96_{-3}^{+5+16}	[Baikov, Chetyrkin, Kühn,'05]
	104 ± 28	[Narison'05]
τ decays \oplus sum rules	99 ± 28	[Narison'05]
Lattice QCD ($n_f = 2$)	97 ± 22	[ALPHA Collab.'05]
	$111 \pm 6 \pm 4 \pm 6$	[QCDSF-UKQCD Collab. '06]
	$101 \pm 8_{-0}^{+25}$	[SPQcdR Collab.'05]
Lattice QCD ($n_f = 3$)	$76 \pm 3 \pm 7$	[HPQCD-MILC-UKQCD Collab.
	86.7 ± 5.9	[CP-PACS-JLQCD Collab. '05]
	$87 \pm 4 \pm 4$	[HPQCD '05]
PDG average	$80 -130$	[Particle Data Group]